# Generalization of a Hardy-Littlewood-Polya Inequality 

Parviz Khajeh-Khalild<br>Christopher Newport College, 50 Shoe Lane, Newport News, Virginia 23606, U.S.A.<br>Communicated by O. Shisha

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#### Abstract

This article is concerned with a generalization of the well-known Hardy-Littlewood-Polya (HLP) inequality to higher dimensions $n \geqslant 2$. We also show via construction of a counterexample that for certain exponents and consequently in some spaces such extension is impossible. © 1991 Academic Press, Inc.


## 1. Introduction

In this paper, we generalize the well-known Hardy-Littlewood-Polya (HLP) inequality (1.1)-(1.3) to higher dimensions. This inequality has applications in many areas of mathematics including approximation of functions; see Theorem 3.4.6 in [1]. Let us recall the classical HLP inequality:

Theorem (G. H. Hardy, J. E. Littlewood, and G. Polya [2, Theorem 330]). Let $1<p<\infty, \varepsilon \neq p-1$. Further, let $f$ be a function defined on $(0, \infty)$ and such that

$$
\int_{0}^{\infty}|f(t)|^{p} t^{\varepsilon} d t<\infty
$$

Then the following inequality holds:

$$
\begin{equation*}
\int_{0}^{\infty}|F(t)|^{p} t^{\varepsilon-p} d t \leqslant\left(\frac{1}{|\varepsilon-p+1|}\right)^{p} \int_{0}^{\infty}|f(t)|^{p} t^{\varepsilon} d t \tag{1.1}
\end{equation*}
$$

where

$$
F(t)=\left\{\begin{array}{lcc}
\int_{0}^{2}|f(s)| d s & \text { for } & \varepsilon<p-1  \tag{1.2}\\
\int_{0}^{\infty}|f(s)| d s & \text { for } & \varepsilon>p-1
\end{array}\right.
$$

Our generalization of (1.1)-(1.3) to a higher dimension, $n \geqslant 2$, is

$$
\begin{equation*}
\int_{\mathfrak{R}^{n}} \frac{|u(x)|^{p}}{|x|^{k+2}} d x \leqslant\left(\frac{p}{|n-k-2|}\right)^{p} \int_{\mathfrak{R}^{n}} \frac{|\nabla u|^{p}}{|x|^{k-p+2}} d x \tag{1.4}
\end{equation*}
$$

where $p$ is a real number $\geqslant 2, x \in \mathfrak{R}^{n}, \nabla=\operatorname{grad},|x|=$ Euclidean norm of $x$ in $\mathfrak{R}^{n}$, and $u$ lies in a proper function space. To show that (1.4) in fact generalizes (1.1)-(1.3) to a higher dimension, choose as $u$ a radial function (i.e., a function of $|x|$ ), $\varepsilon-p=n-k-3$, and let $r=|x|, w_{n}=$ surface area of unit sphere in $\mathfrak{R}^{n}$. Then we have

$$
\begin{align*}
& \int_{\mathfrak{R}^{n}} \frac{|u(x)|^{p}}{|x|^{k+2}} d x=\int_{\mathfrak{R}^{n}} \frac{|u(x)|^{p}}{|x|^{n-\varepsilon+p-1}} d x=\int_{0}^{\infty} \int_{|x|=1} \frac{|u(r)|^{p}}{r^{n-\varepsilon+p-1}} r^{n-1} d w_{n} d r  \tag{1.5}\\
& \int_{\mathfrak{R}^{n}} \frac{|u(x)|^{p}}{|x|^{k+2}} d x=w_{n} \int_{0}^{\infty}|u(r)|^{p} r^{\varepsilon-p} d r .
\end{align*}
$$

Similarly

$$
\begin{align*}
\int_{\mathfrak{R}^{n}} \frac{|\nabla u(x)|^{p}}{|x|^{k-p+2}} d x & =\int_{0}^{\infty} \int_{|x|=1} \frac{|d u(r) / d r|}{r^{n-\varepsilon-1}} r^{n-1} d w_{n} d r \\
& =w_{n} \int_{0}^{\infty}\left|\frac{d u(r)}{d r}\right|^{p} r^{\varepsilon} d r \tag{1.6}
\end{align*}
$$

Substitute (1.5) and (1.6) into (1.4) to see that (1.4) reduces to (1.1). The generalized HLP inequality (1.4) for $p=2$ and $k=0$ has been used in deriving a priori estimates of solutions to some nonlinear partial differential equations (see $[4,6]$ ). One application of (1.4) in particular is to obtain estimates of the solution of semilinear Klein-Gordon equations in $\mathfrak{R}^{3}$

$$
\begin{equation*}
\square u(t, x)+\beta u|u|^{p-1}(t, x)=0 \tag{1.7}
\end{equation*}
$$

where

$$
\square=\frac{\partial^{2}}{\partial t^{2}}-\sum_{i=1}^{3} \frac{\partial^{2}}{\partial x_{i}^{2}} \quad \text { and } \quad x \in \mathfrak{R}^{3}
$$

The estimate obtained then is used to prove the existence of solutions to (1.7). A solution to (1.7) can be expressed in terms of an integral equation (D'Alembert Formula)

$$
\begin{equation*}
u(t, x)=u_{0}(t, x)-\frac{\beta}{4 \sqrt{2 \pi}} \int_{|y-x|=t-\tau} \frac{u|u|^{p-1}(\tau, y)}{|y-x|} d y \tag{1.8}
\end{equation*}
$$

where the integral is taken over the backward light cone with vertex at ( $t, x$ ), and $u_{0}$ is the solution of $\square u_{0}(t, x)=0$. If one uses the change of variable $Y=y-x$ and $d y=\sqrt{2} d Y$, the integral in (1.8) becomes

$$
\begin{equation*}
\frac{\beta \sqrt{2}}{4 \sqrt{2 \pi}} \int_{|Y| \leqslant t-\tau} \frac{u|u|^{p-1}(t-|Y|)}{|Y|} d Y . \tag{1.9}
\end{equation*}
$$

One can use an inequality of type (1.4) to estimate (1.9). As an exampie, consider $p=5$. Then using Holder's inequality, we obtain

$$
\begin{align*}
& \left|\frac{\beta}{4 \sqrt{\pi}} \int_{|Y| \leqslant t-\tau} \frac{u^{5}(t-|Y|, Y)}{|Y|} d Y\right| \\
& \quad \leqslant \frac{\beta}{4 \sqrt{\pi}}\|u\|_{L^{\infty}}\left[\int_{|Y| \leqslant t-\tau} u^{6}(t-|Y|, Y) d Y\right]^{1 / 2} \\
& \quad \times\left[\int_{|Y| \leqslant t-\tau} \frac{u^{2}(t-|Y|, Y)}{|Y|^{2}} d Y\right]^{1 / 2} \tag{1.10}
\end{align*}
$$

The first integral on the right can be approximated via an energy estimate, while the second one is of type (1.4) for $n=3, p=2$, and $k=0$, but this can also be approximated by an energy estimate. For details and further discussion of this problem, see [6]. An estimate of the type (1.10) that leads to existence and uniqueness of the solution to the semilinear Klein-Gordon equation (1.7) motivates the search for a more general inequality

$$
\begin{equation*}
\left[\int_{\mathfrak{F}^{n}} \frac{|u(x)|^{p}}{|x|^{l}} d x\right]^{1 / p} \leqslant C\left[\int_{\mathfrak{F}^{n}} \frac{|\nabla u(x)|^{q}}{|x|^{s}} d x\right]^{1 / q} \tag{1.11}
\end{equation*}
$$

where $p, q, C, s$, and $l$ are positive real numbers. Theorems 1 and 2 show that (1.11) holds for $p=q \geqslant 2$,

$$
l-s=p, \quad C=\frac{p}{|n-l|},
$$

and $u \in C_{0}^{\infty}\left(\Re^{n} \backslash\{0\}\right)$, the set of all infinitely differentiable functions with compact support in $\mathfrak{R}^{n} \backslash\{0\}, n \geqslant 2$. For other choices of $p, q, l$, and $C$ the question remains open. In Section 3 we verify, by using a counterexample, that when $l=n$, inequality (1.4) does not hold.

## 2. Statement of the Result

Let $n \in N$ and $k \in \mathfrak{R}$ and assume $n-k \neq 2$; then for all $u \in C_{0}^{\infty}\left(\mathfrak{R}^{n} \backslash\{0\}\right)$

$$
\begin{equation*}
\int_{\mathfrak{R}^{n}} \frac{u^{2}(x)}{|x|^{k+2}} d x \leqslant \frac{4}{(n-k-2)^{2}} \int_{\Re^{n}} \frac{|\nabla u|^{2}}{|x|^{k}} d x . \tag{2.1}
\end{equation*}
$$

For $2 \leqslant p<\infty, n-k \neq 2$, and for all $u \in C_{0}^{\infty}\left(\Re^{n} \backslash\{0\}\right)$ we have the following generalization of (2.1):

$$
\begin{equation*}
\int_{\mathfrak{R}^{n}} \frac{|u(x)|^{p}}{|x|^{k+2}} d x \leqslant\left(\frac{p}{|n-k-2|}\right)^{p} \int_{\mathfrak{R}^{n}} \frac{|\nabla u|^{p}}{|x|^{k-p+2}} d x \tag{2.2}
\end{equation*}
$$

This inequality will be referred to as a "generalized" Hardy-LittlewoodPolya inequality since it reduces to (1.1) when $u$ is a radial function. This is obviously a generalization of the HLP inequality only when $p \geqslant 2$. We summarize all these in Theorem 1 and Theorem 2.

Theorem 1. Let $n \in N, k \in \Re$ and $n-k \neq 2$. Then for all $u \in C_{0}^{\infty}\left(\Re^{n} \backslash\{0\}\right)$ and $2 \leqslant p<\infty$, the inequality (2.1) holds.

Proof. We search for an identity,

$$
\begin{equation*}
\int_{\mathfrak{R}^{n}}\left(\frac{|\nabla u|^{2}}{r^{k}}-\frac{1}{C^{2}} \frac{u^{2}}{r^{k+2}}\right) d x=\int_{\mathfrak{R}^{n}} \frac{1}{h(r)}|\nabla(g(r) u)|^{2} d x \tag{2.3}
\end{equation*}
$$

where $r=|x|$ and $g, h$ are two positive functions to be determined later, along with the constant $C$. Let

$$
\Lambda=\frac{1}{h(r)}|\nabla(g(r) u)|^{2}
$$

where $u=u(x)$; then

$$
\begin{aligned}
A & =\frac{1}{h(r)}\left|g^{\prime}(r) u \frac{x}{r}+g(r) \nabla u\right|^{2} \\
& =\frac{1}{h(r)}\left\{2 g^{\prime}(r) g(r) u\left(\frac{x}{r} \cdot \nabla u\right)+g^{\prime}(r)^{2} u^{2}+g^{2}(r)|\nabla u|^{2}\right\} .
\end{aligned}
$$

Now consider

$$
\begin{align*}
\int_{\mathfrak{R}^{n}} \Lambda d x= & \int_{\mathfrak{R}^{n}} \frac{2 g^{\prime}(r) g(r)}{h(r)} u\left(\frac{x}{r} \cdot \nabla u\right) d x \\
& +\int_{\mathfrak{R}^{n}} \frac{g^{\prime}(r)^{2}}{h(r)} u^{2} d x+\int_{\mathbb{R}^{n}} \frac{g^{2}(r)}{h(r)}|\nabla u|^{2} d x \tag{2.4}
\end{align*}
$$

The first integral on the right hand side can be simplified. After using Green's identity we obtain

$$
\begin{align*}
\int_{\mathfrak{R}^{n}} \frac{2 g^{\prime}(r) g(r)}{r h(r)} u\left(\sum_{i=1}^{n} x_{i} u_{x_{i}}(x)\right) d x & =\sum_{i=1}^{n} \int_{\mathfrak{R}^{n}} \frac{g^{\prime}(r) g(r)}{r h(r)} x_{i}\left(\frac{\partial u^{2}}{\partial x_{i}}\right) d x \\
& =-\int_{\mathfrak{R}^{n}} u^{2}(x) \operatorname{div}\left(\frac{g^{\prime}(r) g(r) x}{r h(r)}\right) d x . \tag{2.5}
\end{align*}
$$

Use (2.4), (2.5) to expand (2.3) as

$$
\begin{align*}
& \int_{\mathfrak{R}^{n}}\left(\frac{|\nabla u|^{2}}{r^{k}}-\frac{1}{C^{2}} \frac{u^{2}}{r^{k+2}}\right) d x \\
& \quad=-\int_{\mathfrak{R}^{n}} u^{2}(x) \operatorname{div}\left(\frac{g^{\prime}(r) g(r) x}{r h(r)}\right) d x+\int_{\mathfrak{R}^{n}} \frac{g^{\prime}(r)^{2} u^{2}}{h(r)} d x \\
& \quad+\int_{\mathfrak{R}^{n}} \frac{g^{2}(r)}{h(r)}|\nabla u|^{2} d x . \tag{2.6}
\end{align*}
$$

The above identity holds provided

$$
\begin{equation*}
\frac{1}{r^{k}}=\frac{g^{2}(r)}{h(r)} \tag{2.7}
\end{equation*}
$$

and

$$
\begin{equation*}
\frac{1}{C^{2} r^{k+2}}=\operatorname{div}\left(\frac{g^{\prime}(r) g(r) x}{r h(r)}\right)-\frac{g^{\prime}(r)^{2}}{h(r)} . \tag{2.8}
\end{equation*}
$$

Solving (2.7) for $h$ and then substituting into (2.8) yields

$$
\begin{equation*}
\frac{1}{C^{2} r^{k+2}}=\operatorname{div}\left(\frac{g^{\prime}(r) x}{r^{k+1} g(r)}\right)-\frac{g^{\prime}(r)^{2}}{r^{k} g^{2}(r)} \tag{2.9}
\end{equation*}
$$

Simplify (2.9) as

$$
\frac{1}{C^{2}} g^{2}=n r g g^{\prime}+r^{2} g g^{\prime \prime}-(k+1) r g g^{\prime}-2 r^{2} g^{\prime 2}
$$

or

$$
\begin{equation*}
r^{2} g g^{\prime \prime}+(n-k-1) r g g^{\prime}-2 r^{2} g^{\prime 2}-\frac{1}{C^{2}} g^{2}=0 \tag{2.10}
\end{equation*}
$$

This nonlinear O.D.E. in $g$ has a solution of the form $g(r)=r^{s}$. Substituting $g(r)=r^{s}$ in (2.10), we see that $s$ is a solution of the equation

$$
s^{2}-(n-k-2) s+\frac{1}{C^{2}}=0 .
$$

The roots are real (which is needed to fulfill (2.5)) provided that

$$
(n-k-2)^{2}-\frac{4}{C^{2}} \geqslant 0
$$

Therefore

$$
\frac{1}{C^{2}} \leqslant \frac{(n-k-2)^{2}}{4} \quad \text { or } \quad C^{2} \geqslant \frac{4}{(n-k-2)^{2}}
$$

And, for $C^{2}=(4) /(n-k-2)$, there is a double root as $s=(n-k-2) / 2$.

$$
\text { Thus }\left\{\begin{align*}
g(r) & =r^{(n-k-2) / 2}  \tag{2.11}\\
h(r) & =r^{n-2} \\
C^{2} & =\frac{4}{(n-k-2)^{2}}
\end{align*}\right.
$$

This proves that (2.3) is valid when $g, h$, and $C$ satisfy (2.11). Inequality (2.1) follows immediately.

Theorem 2. Let $n \in N, k \in \mathfrak{R}$ and $n-k \neq 2$. Then for all $u \in C_{0}^{\infty}\left(\mathfrak{R}^{n} \backslash\{0\}\right)$ and $2 \leqslant p<\infty$, the inequality (2.2) holds.

Proof. For $p=2$ the inequality holds because of Theorem 1. Now let $p>2$ and $\varepsilon>0$ for $u \in C_{0}^{\infty}\left(\Re^{n} \backslash\{0\}\right)$, and let $v(x)=\left|u^{2}(x)+\varepsilon\right|^{p / 4} \chi(x)$, where $\chi(x) \equiv 1$ on the support of $u$ and $\chi \in C_{0}^{\infty}\left(\Re^{n} \backslash\{0\}\right)$. Thus $v \in C_{0}^{\infty}\left(\mathfrak{R}^{n} \backslash\{0\}\right)$ so that it satisfies (2.1) and

$$
\begin{equation*}
\int_{\mathfrak{R}^{n}} \frac{v^{2}(x)}{|x|^{k+2}} d x \leqslant \frac{4}{(n-k-2)^{2}} \int_{\mathfrak{R}^{n}} \frac{|\nabla v(x)|^{2}}{|x|^{k}} d x \tag{2.12}
\end{equation*}
$$

as well.
Expanding

$$
|\nabla v(x)|^{2}=\left|\frac{p}{4}\right| u^{2}+\left.\varepsilon\right|^{(p / 4)-1} 2 u(\nabla u) \chi(x)+\left.\left|u^{2}+\varepsilon\right|^{p / 4} \nabla \chi\right|^{2}
$$

and using the orthogonality of the vectors $\nabla u$ and $\nabla \chi, \nabla u \cdot \nabla \chi=0$, we arrive at

$$
|\nabla v(x)|^{2}=\frac{p^{2}}{4}\left|u^{2}+\varepsilon\right|^{(p / 2)-2} u^{2} \chi^{2}|\nabla u|^{2}+\left|u^{2}+\varepsilon\right|^{p / 2}|\nabla \chi|^{2}
$$

Therefore

$$
\int_{\mathfrak{R}^{n}} \frac{|\nabla v(x)|^{2}}{|x|^{k}} d x=\frac{p^{2}}{4} \int_{\mathfrak{R}^{n}} \frac{\left|/ u^{2}+\varepsilon\right|^{p / 2-2} u^{2}|\nabla u|^{2}}{|x|^{k}} d x+\int_{\mathfrak{R}^{n}} \frac{\left|u^{2}+\varepsilon\right|^{p / 2}}{|x|^{k}}|\nabla \chi|^{2} d x .
$$

Now let $\varepsilon \rightarrow 0$,

$$
\begin{equation*}
\lim _{\varepsilon \rightarrow 0} \int_{\mathfrak{R}^{n}} \frac{|\nabla v(x)|^{2}}{|x|^{k}} d x=\frac{p^{2}}{4} \int_{\mathfrak{R}^{n}} \frac{|u|^{p-2}|\nabla u|^{2}}{|x|^{k}} d x \tag{2.13}
\end{equation*}
$$

and

$$
\begin{equation*}
\lim _{\varepsilon \rightarrow 0} \int_{\mathfrak{R}^{n}} \frac{v^{2}(x)}{|x|^{k+2}} d x=\int_{\mathfrak{R}^{n}} \frac{|u|^{p}}{|x|^{k+2}} d x \tag{2.14}
\end{equation*}
$$

Substituting in (2.12) yields

$$
\begin{equation*}
\int_{\mathfrak{R}^{n}} \frac{|u|^{p}}{|x|^{k+2}} d x \leqslant \frac{4}{(n-k-2)^{2}}\left(\frac{p^{2}}{4}\right) \int_{\mathfrak{P}^{n}} \frac{|u|^{p}{ }^{2}|\nabla u|^{2}}{|x|^{k}} d x . \tag{2.15}
\end{equation*}
$$

We use Holder's inequality on the right-hand side to show that

$$
\begin{equation*}
\int_{\mathfrak{R}^{n}} \frac{|u|^{p-2}|\nabla u|^{2}}{|x|^{k}} d x \leqslant\left(\int_{\mathfrak{R}^{n}} \frac{|u|^{q(p-2)}}{|x|^{k_{1} q}} d x\right)^{1 / q}\left(\int_{\mathfrak{R}^{n}} \frac{|\nabla u|^{2 t}}{|x|^{k_{2} t}} d x\right)^{1 / t}, \tag{2.16}
\end{equation*}
$$

where $t=p / 2,1 / q+1 / t=1$, and $k_{1}+k_{2}=k$. Now let $k_{1} q=k+2$. Then $k_{1}=(k+2) / q$ and $k_{2}=k-(k+2) / q$. Also $q(p-2)=p$, since $1 / q=1-2 / p$. Therefore, it is easy to establish that $k_{2} t=k-p+2$. Substitution into (2.16) gives

$$
\begin{equation*}
\int_{\mathfrak{R}^{n}} \frac{|u|^{p-2}|\nabla u|^{2}}{|x|^{k}} d x \leqslant\left(\int_{\mathfrak{R}^{n}} \frac{|u|^{p}}{|x|^{k+2}} d x\right)^{1 / q}\left(\int_{\mathfrak{R}^{n}} \frac{|\nabla u|^{p}}{|x|^{k-p-2}} d x\right)^{2 / p}, \tag{2.17}
\end{equation*}
$$

which is in turn substituted in (2.15) as

$$
\int_{\mathfrak{R}^{n}} \frac{|u|^{p}}{|x|^{k+2}} d x \leqslant \frac{p^{2}}{(n-k-2)^{2}}\left(\int_{\mathfrak{R}^{n}} \frac{|u|^{p}}{|x|^{k+2}} d x\right)^{1 / q}\left(\int_{\mathfrak{R}^{n}} \frac{|\nabla u|^{p}}{|x|^{k-p+2}} d x\right)^{2 / p}
$$

or

$$
\left(\int_{\mathfrak{R}^{n}} \frac{|u|^{p}}{|x|^{k+2}} d x\right)^{1-1 / q} \leqslant \frac{p^{2}}{(n-k-2)^{2}}\left(\int_{\mathfrak{R}^{n}} \frac{|\nabla u|^{p}}{|x|^{k-p+2}} d x\right)^{2 / p}
$$

But $1-1 / q=1 / t=2 / p$ so that we have

$$
\left(\int_{\Re^{n}} \frac{|u|^{p}}{|x|^{k+2}} d x\right)^{2 / p} \leqslant \frac{p^{2}}{(n-k-2)^{2}}\left(\int_{\mathfrak{R}^{n}} \frac{|\nabla u|^{p}}{|x|^{k-p+2}} d x\right)^{2 / p}
$$

which proves (2.2).

## 3. Construction of a Counterexample

Here we show that when $n=k+2$, there is a sequence $u_{m} \in C_{0}^{\infty}\left(\Re^{n} \backslash\{0\}\right)$ such that as $m \rightarrow \infty$,

$$
\int_{\mathfrak{M}^{n}} \frac{u_{m}^{2}(x)}{|x|^{k+2}} d x \rightarrow \infty,
$$

while

$$
\int_{\mathfrak{M} n} \frac{\left|\nabla u_{m}(x)\right|^{2}}{|x|^{k}} d x<\infty .
$$

This shows that when $n=k+2$, an inequality similar to (2.1) does not exist. We consider $u_{m}(x)$ to be a radial function; that is, $u_{m}(x)=u_{m}(r)$. Then

$$
\int_{\mathfrak{M}^{n}} \frac{u_{m}^{2}(x)}{|x|^{n}} d x=w_{n} \int_{0}^{\infty} \frac{u_{m}^{2}(r)}{r} d r,
$$

where $w_{n}=$ surface area of unit sphere in $\mathfrak{R}^{n}$. Similarly,

$$
\int_{\mathfrak{R}^{n}} \frac{\left|\nabla u_{m}(x)\right|^{2}}{|x|^{n-2}} d x=w_{n} \int_{0}^{\infty} r u_{m}^{\prime}(r)^{2} d r
$$

The heart of the construction is the fact that the function

$$
g(x)=\frac{1}{\sqrt{\ln x}}
$$

satisfies

$$
\int_{2}^{\infty} \frac{g(x)^{2}}{x} d x=+\infty
$$

and

$$
\int_{2}^{\infty} x g^{\prime}(x)^{2} d x<\infty
$$

Choose $f \geqslant 0, f \in C_{0}^{\infty}(\mathfrak{R})$, so that $f=0$ on $(-\infty, 1]$ and

$$
f(x)=\frac{1}{\sqrt{\ln x}} \quad \text { for } \quad x \geqslant 2 .
$$

Then

$$
\begin{gathered}
\int_{0}^{\infty} \frac{f(x)^{2}}{x} d x=+\infty \\
\int_{0}^{\infty} x f^{\prime}(x)^{2} d x<\infty
\end{gathered}
$$

Also, choose $\Psi \in C_{0}^{\infty}(\mathfrak{R})$ with $\operatorname{supp} \Psi=[-2,2], \Psi$ equal to one on $[-1,1]$, and $0 \leqslant \Psi \leqslant 1$ elsewhere.

Let $u_{m}(x)=\Psi(x / m) f(x)$. Then

$$
\int_{0}^{\infty} \frac{u_{m}^{2}(x)}{x} d x \geqslant \int_{2}^{m} \frac{f(x)^{2}}{x} d x \rightarrow \infty \quad \text { as } \quad m \rightarrow+\infty
$$

On the other hand

$$
\begin{aligned}
u_{m}^{\prime}(x) & =\frac{1}{m} \Psi^{\prime}\left(\frac{x}{m}\right) f(x)+\Psi\left(\frac{x}{m}\right) f^{\prime}(x) \\
\left|u_{m}^{\prime}(x)\right| & \leqslant \frac{1}{m}\left|\Psi^{\prime}\right|_{L^{\infty}(\mathfrak{R})}+\left|f^{\prime}(x)\right|
\end{aligned}
$$

or

$$
\left|u_{m}^{\prime}(x)\right| \leqslant \frac{C}{m}+\left|f^{\prime}(x)\right|
$$

Using the inequality $a b \leqslant \frac{1}{2}\left(a^{2}+b^{2}\right)$ we find that

$$
u_{m}^{\prime}(x)^{2} \leqslant D\left(\frac{1}{m^{2}}+f^{\prime 2}\right)
$$

Then since $u_{m}$ vanishes for $n \geqslant 2 m$, we have

$$
\int_{0}^{\infty} x u_{m}^{\prime}(x)^{2} d x \leqslant \int_{0}^{2 m} \frac{D}{m^{2}} x d x+D \int_{0}^{\infty} x f^{\prime}(x)^{2} d x<\infty
$$

The construction is complete.
Remark 1. From the construction of identity (2.3) it is clear that

$$
C=\frac{4}{(n-k-2)^{2}}
$$

is the best possible constant for (2.1) and consequently

$$
C=\left(\frac{p}{|n-k-2|}\right)^{p}
$$

is the best for (2.2). Note that this is consistent with the best possible constant of the Hardy-Littlewood-Polya inequality.

Remark 2. The HLP inequality has an application in weighted Sobolev spaces. See for example [3, p. 28]. The generalized HLP inequality has applications, as we showed in our introduction, in establishing the existence of a solution to nonlinear Klein-Gordon equations (see [6]) and in elliptic equations (see [4, p. 451]).

Remark 3. The inequality (2.2) holds only when $2 \leqslant p<\infty$. The question of whether this can be extended to $1<p<\infty$ is an open question.

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